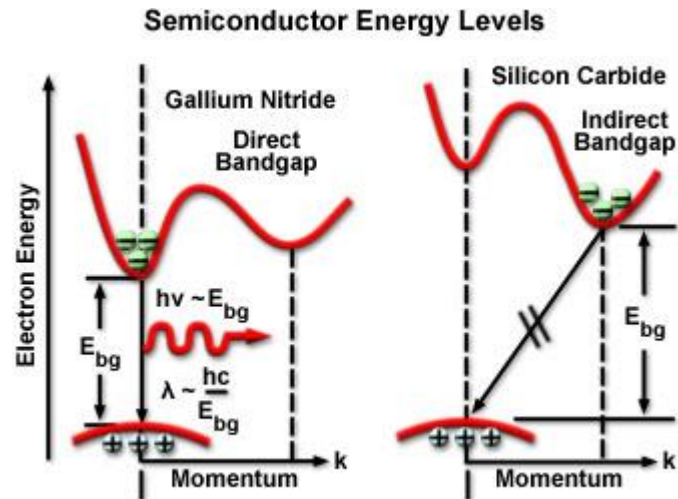


# Spin flip lasers

Tim Verhagen

# Spin flip laser

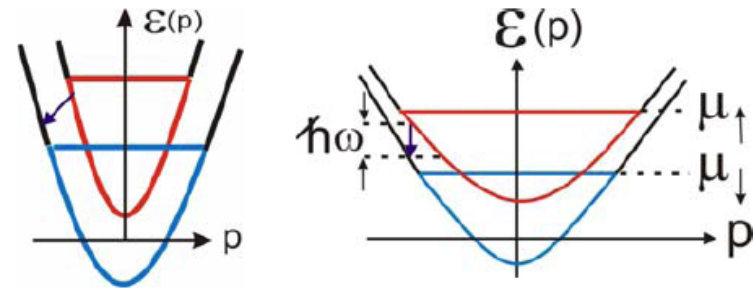
## Semiconductor (laser diode)



$\Delta k = 0$   
photon

$\Delta k \neq 0$   
phonon

## Magnetic (spin flip laser)



$\Delta k \neq 0$

magnon

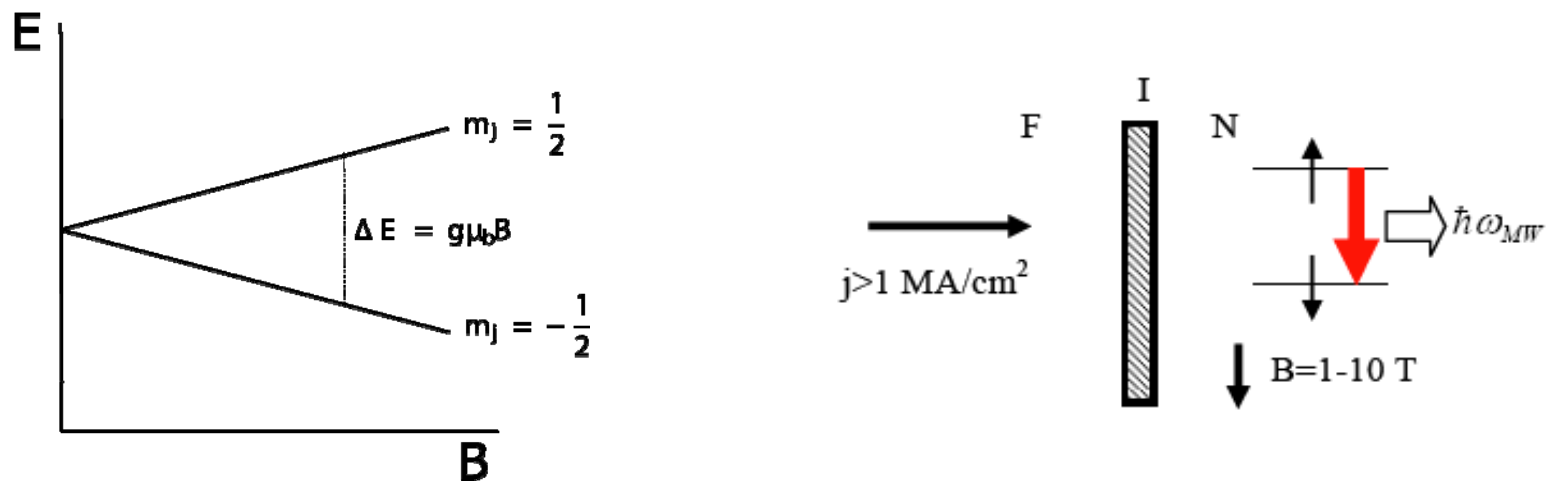
High frequency  
oscillators

$\Delta k = 0$

photon

High frequency  
lasers

# Zeeman split transition



Zeeman effect

$$h\nu = g\mu_B m_s B$$

$$\nu = 0.007gB[\text{THz}]$$

# Devices, Majority F

If coercive field of F is (much) bigger than applied field

(  $\text{SmCo}_5$ ,  $\text{AlNiCo}$ ,  $\text{Nd}_2\text{Fe}_{14}\text{B}$  )

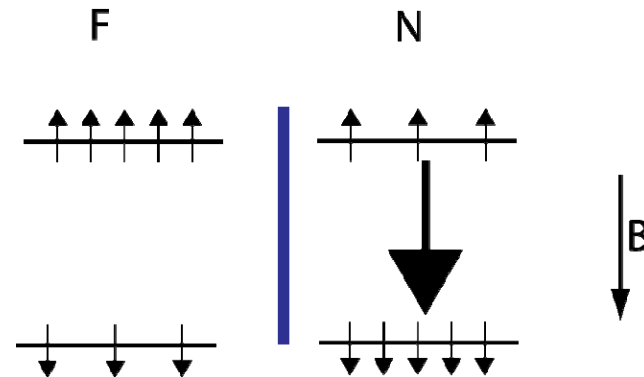
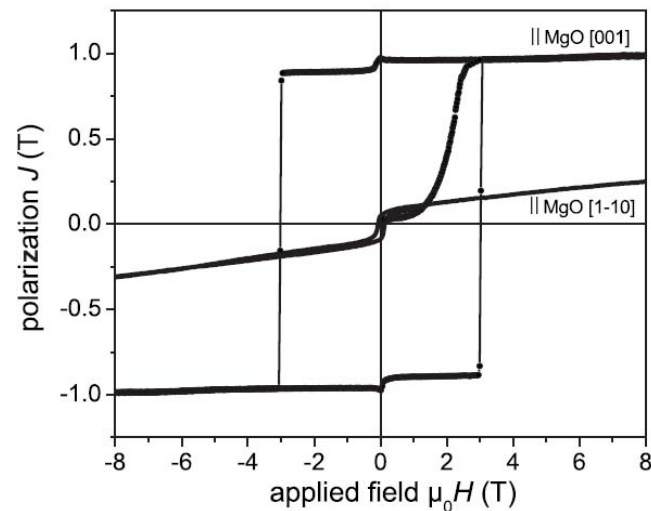
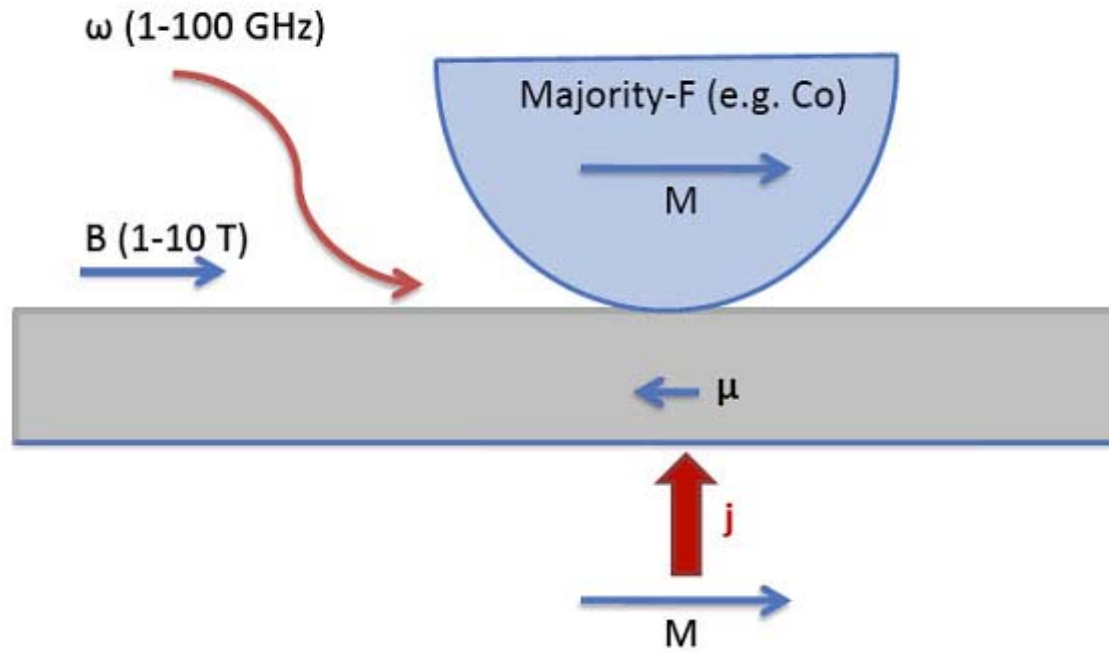


FIG. 2. Magnetic hysteresis of a  $\text{SmCo}_5$  film measured along the easy magnetization axis ( $\parallel \text{MgO} [001]$ ) and along the in-plane hard axis ( $\parallel \text{MgO} [1-10]$ ).

PRB 77 104443

# Experimental Zeeman lasing

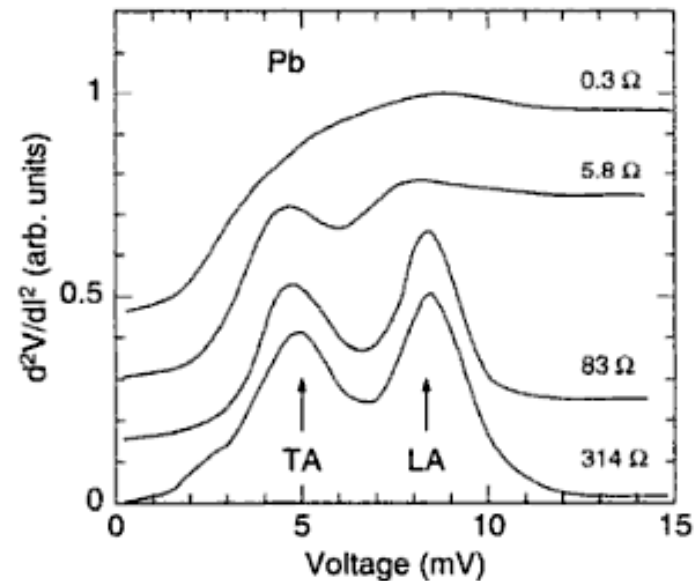


- Spin flip laser
- **Point contact spectroscopy**
- SmCo<sub>5</sub>
- Radiation setup

# What is a good point contact?

1. For pure metals, the contact resistance is within 1 to 100  $\Omega$  and the short has the metallic conductivity; i. e., its resistance increases with the voltage rise so that  $d^2V/dI^2 > 0$  over the whole energy range.
2. Distinct maxima are observed in the PC spectrum at energy between zero and maximal one  $\hbar\omega_{\max}$ . They are reproducible for different contacts of the particular metal. At energy  $eV > \hbar\omega_{\max}$ , the spectrum contains a certain constant background value. The background parameter  $\gamma = B/A$  (Fig. 4.2) is different for each metal, but as a rule, it should not be greater than 0.5.
3. Relative variation of the differential resistance within the spectrum has to be maximal among the spectra with the same other characteristics. Around zero-bias at  $eV \ll \hbar\omega_{\max}$ , the  $d^2V/dI^2$  curve behaves monotonously as  $V^n$  (see Table 3.1) or zero-bias anomaly should be as weak as possible.

# What is a good point contact?

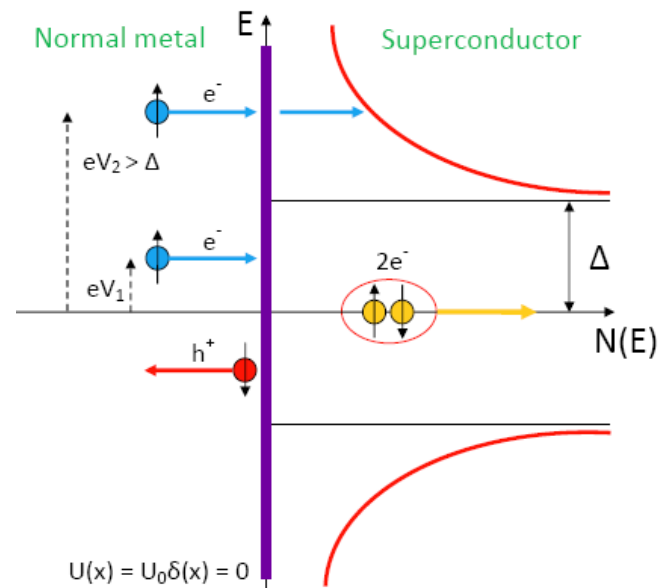


**Fig. 2.3.** Second derivative of  $I - V$  curves for a Pb tunnel junction with micro-constriction at helium temperature in the process of successive decreasing of the resistance of the constriction from 314  $\Omega$  (bottom curve) to 0.3  $\Omega$  (top curve). The positions of longitudinal acoustic (LA) and transverse acoustic (TA) phonon peaks in Pb are indicated by arrows. The curves are offset vertically for clarity. Data taken from Yanson (1974).



# Add a superconductor

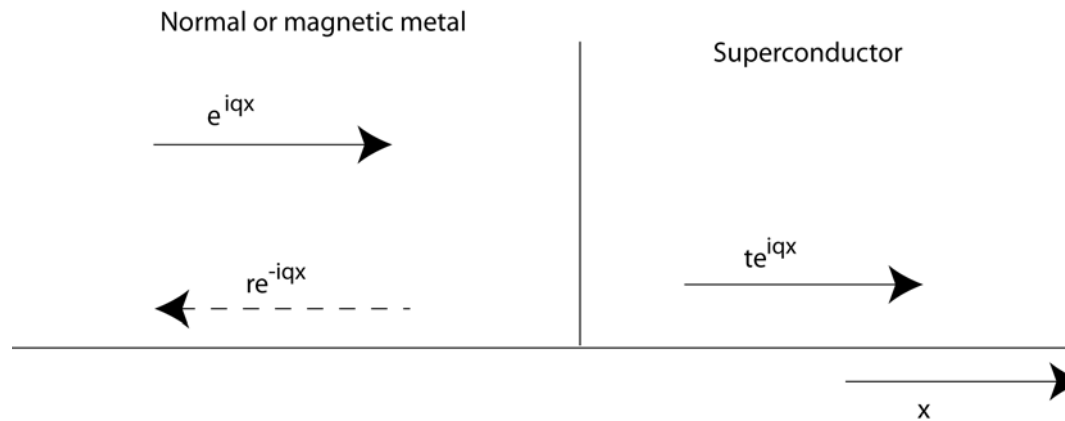
Create a normal metal – superconductor interface,  
so Andreev reflection



# Add a superconductor

Bogoliubov de Gennes equation

$$\begin{cases} i\hbar \frac{\partial f(x,t)}{\partial t} = \left[ -\frac{\hbar^2 \nabla^2}{2m} - \mu(x) + V(x) \right] f(x,t) + \Delta(x)g(x,t) \\ i\hbar \frac{\partial g(x,t)}{\partial t} = \left[ \frac{\hbar^2 \nabla^2}{2m} - \mu(x) + V(x) \right] g(x,t) + \Delta(x)f(x,t) \end{cases}$$

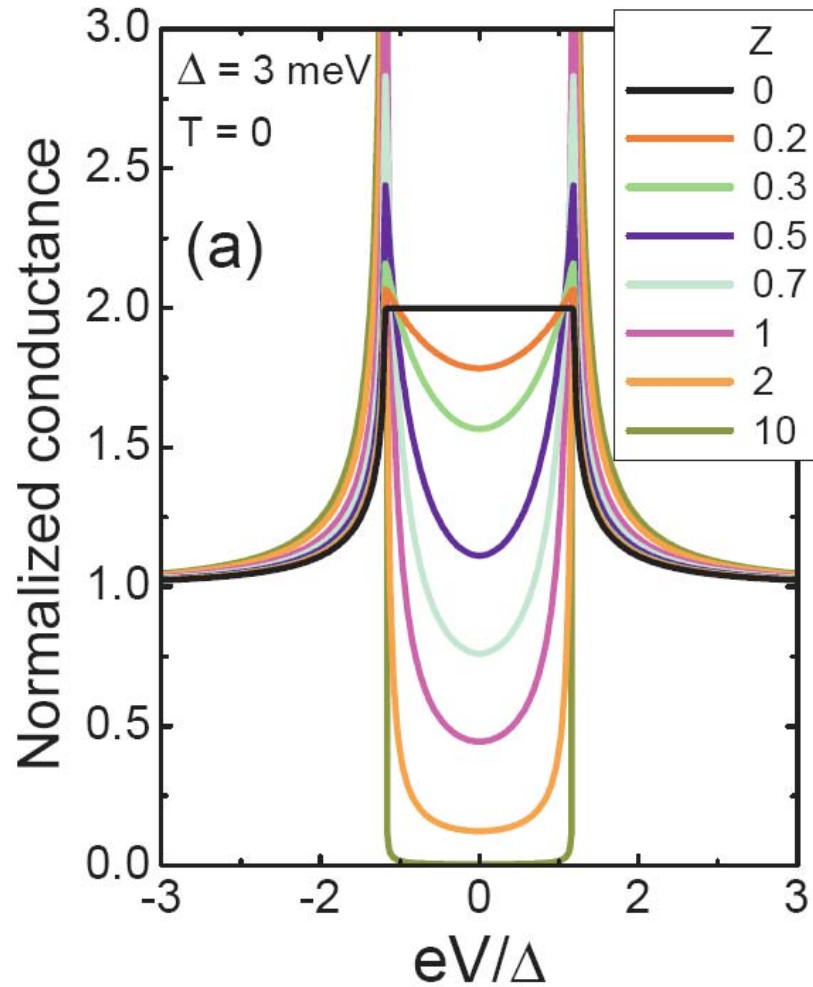


Solve -> Transmission coefficient @ 0K:  $T(\varepsilon) = 1 + A(\varepsilon) - B(\varepsilon)$

Probability Andreev  
reflection

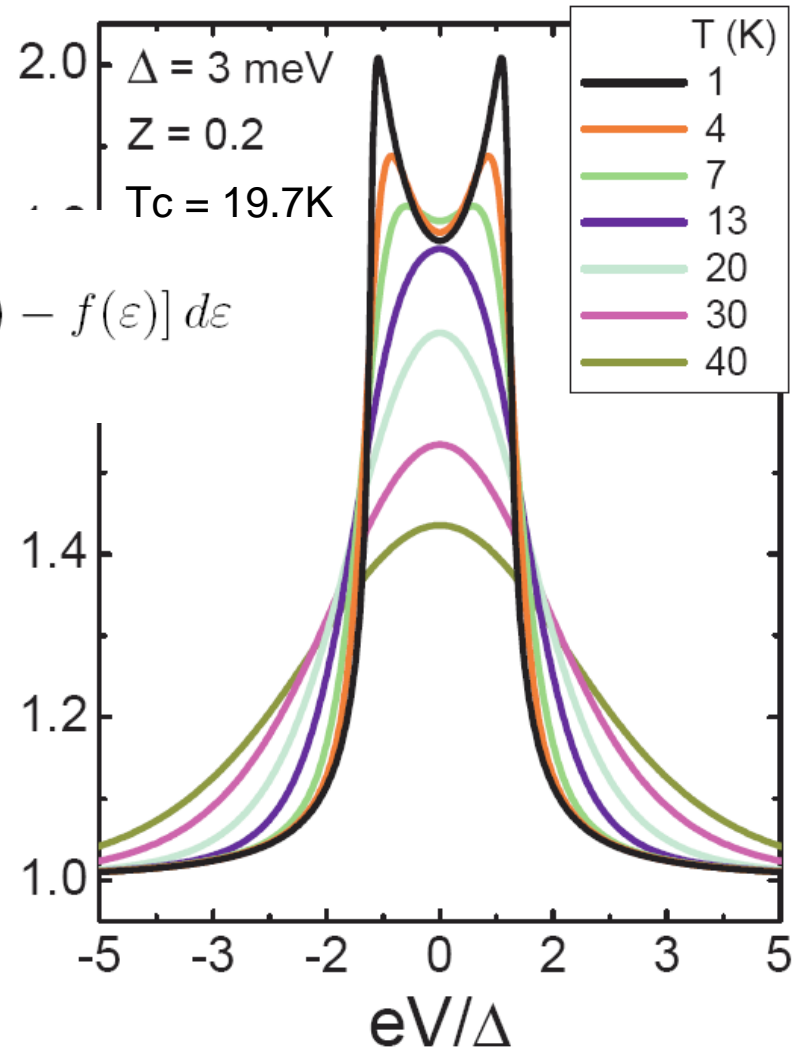
Probability normal  
reflection

# Add a superconductor

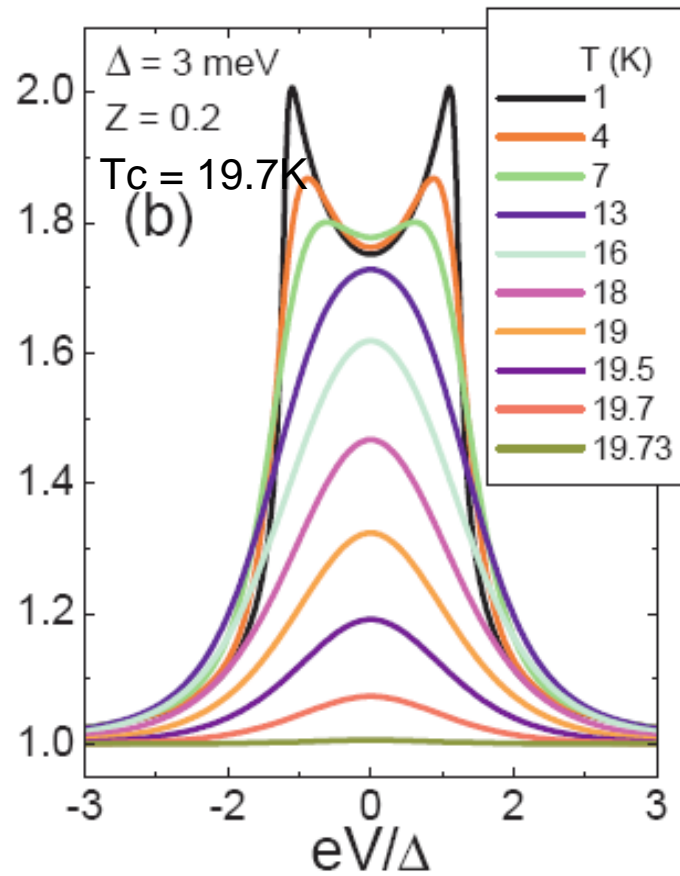


# Add temperature

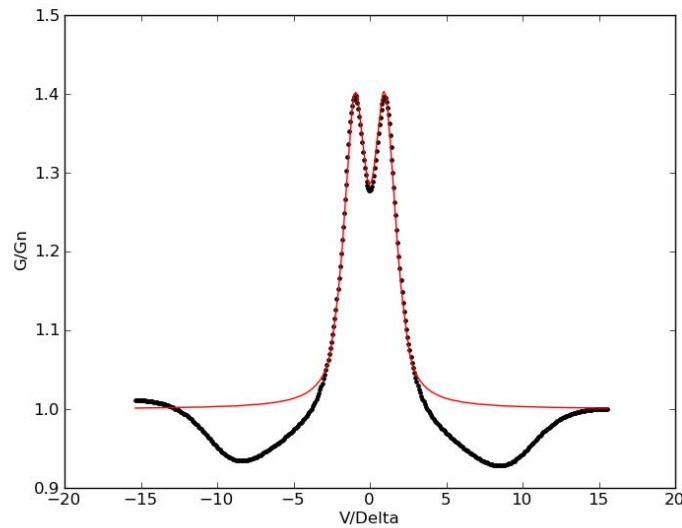
$$I = \int_{-\infty}^{+\infty} [1 + A(\varepsilon) - B(\varepsilon)] [f(\varepsilon - eV) - f(\varepsilon)] d\varepsilon$$



# Include gap(T)



# Too low temperature?



Best  $T < 4.2$  K

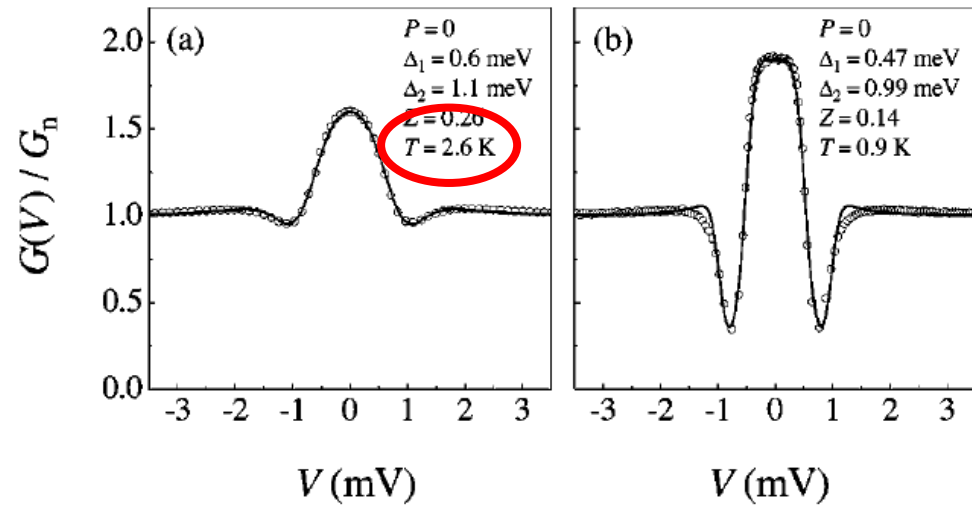



FIG. 4. Representative conductance versus voltage curves for Cu/Nb at  $T = 4.2$  K (open circles) with a contact resistance of (a)  $R = 7.6 \Omega$  and (b)  $R = 2.4 \Omega$ . The solid lines are fits using the model of Sec. II, resulting in  $P$ ,  $\Delta_1$ ,  $\Delta_2$ ,  $Z$ , and  $T$  as indicated in the figure.

the proximity effect in the Cu plays an important role. Since temperature is included in our model by an equilibrium Fermi-Dirac distribution function, the fitted temperature is considerably lower than 4.2 K, as shown in Fig. 4. Formally

Strijkers, PRB **63** 104510

# Add inelastic scattering

$$\begin{aligned}
 i\hbar \frac{\partial f(x,t)}{\partial t} &= \left[ -\frac{\hbar^2 \nabla^2}{2m} - \mu(x) - i\Gamma_i + V(x) \right] f(x,t) + \Delta(x)g(x,t) \\
 i\hbar \frac{\partial g(x,t)}{\partial t} &= \left[ \frac{\hbar^2 \nabla^2}{2m} - \mu(x) + i\Gamma_i + V(x) \right] g(x,t) + \Delta(x)f(x,t)
 \end{aligned}$$


 Inelastic scattering

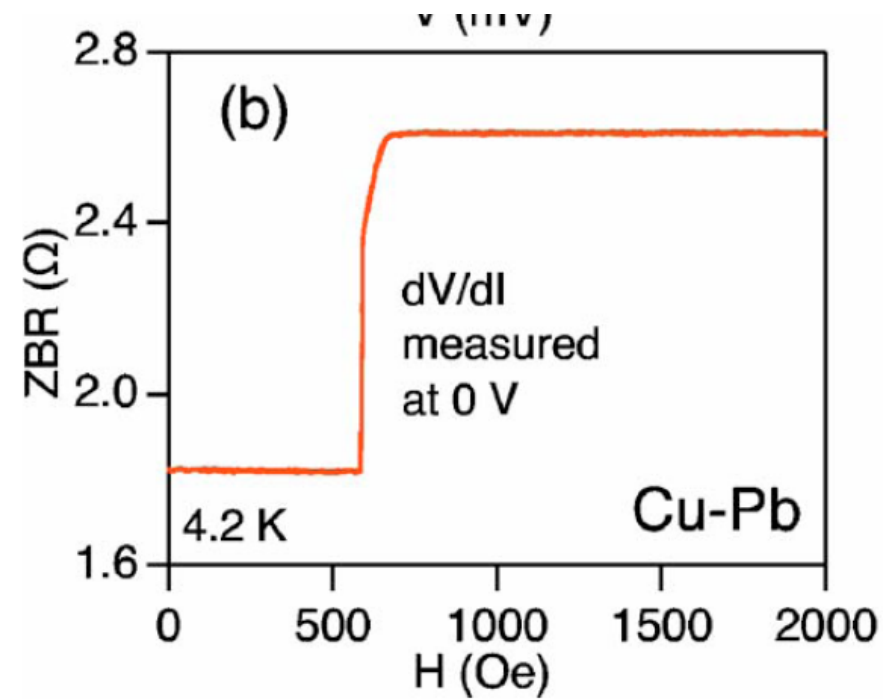
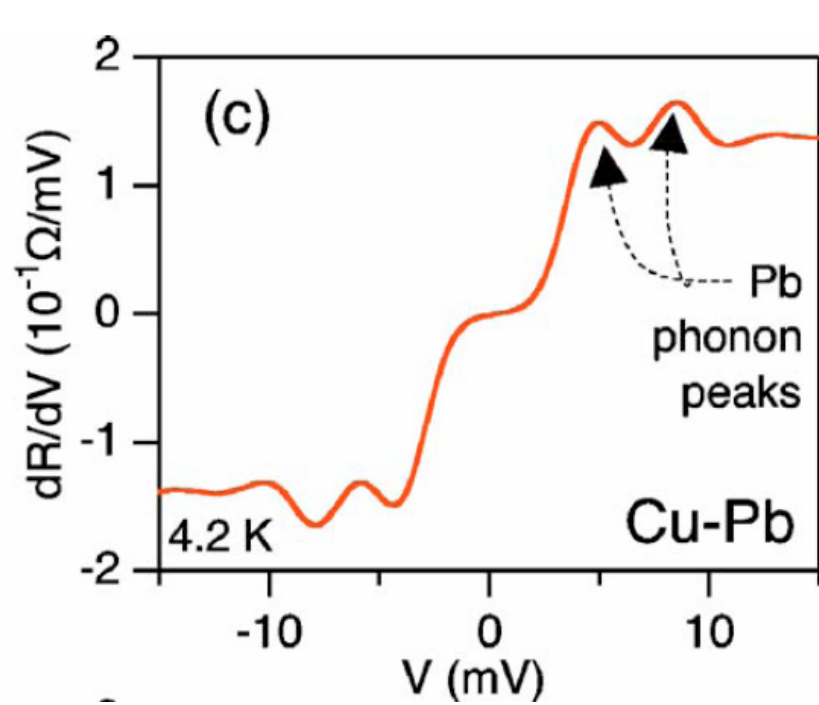
$$u_0^2 = \frac{1}{2} \left[ 1 + \frac{\sqrt{\varepsilon^2 + \Delta^2}}{\varepsilon} \right] \xrightarrow{\text{Dyson}} u_0^2 = \frac{1}{2} \left[ 1 + \frac{\sqrt{(\varepsilon + i\Gamma)^2 + \Delta^2}}{\varepsilon + i\Gamma} \right]$$

Solve -> Transmission coefficient @ 0K:  $T(\varepsilon) = 1 + A(\varepsilon) - B(\varepsilon)$

Probability Andreev  
reflection

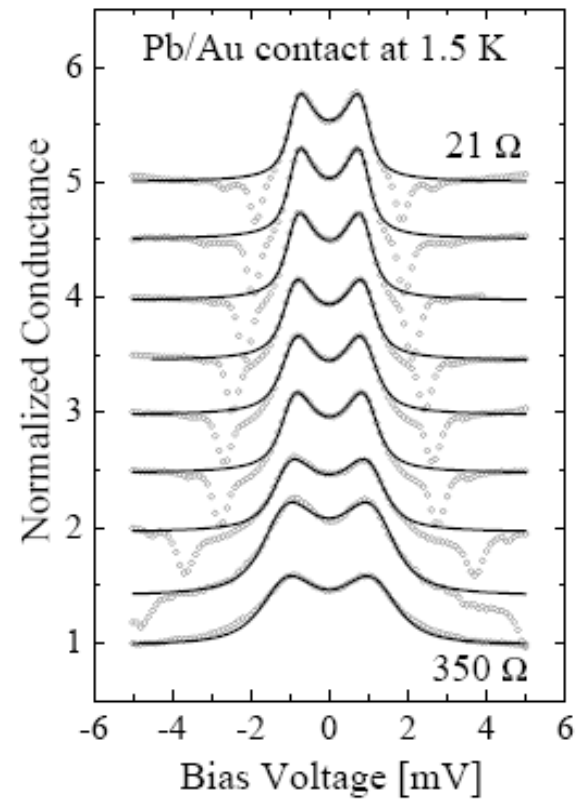
Probability normal  
reflection

# What to expect?

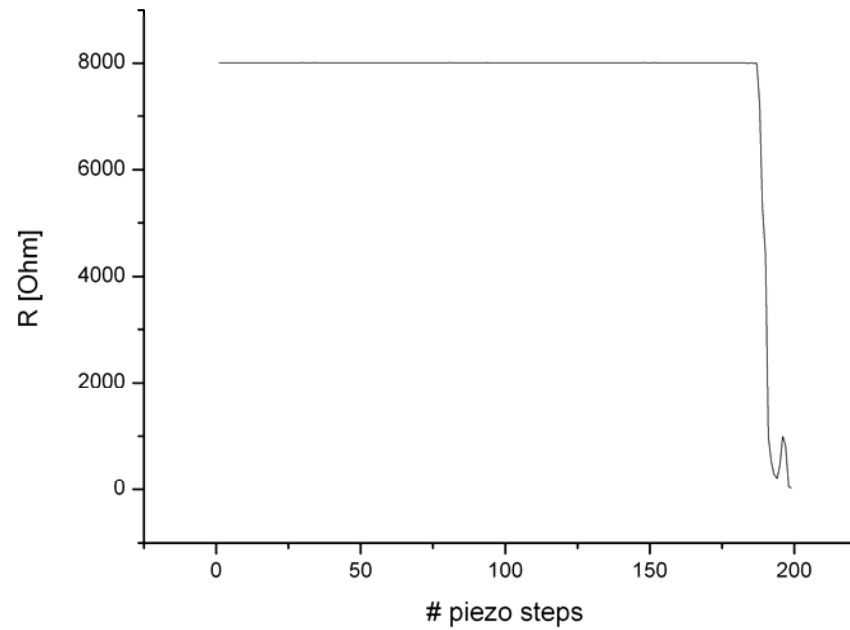




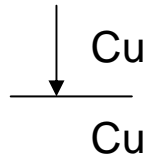
# What to expect?



# Making a point contact with attocube



Together with A.Naylor, Leeds



# Phonon peaks?

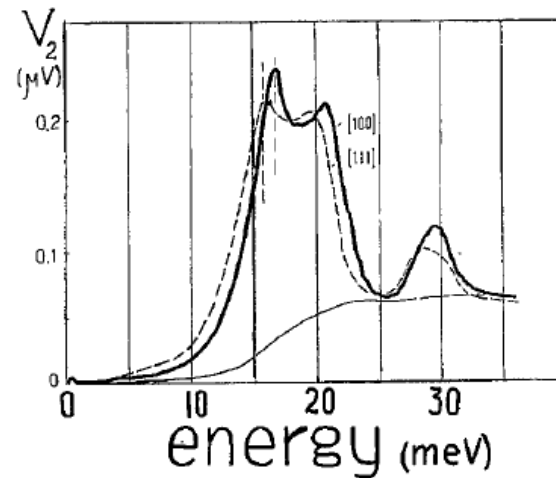
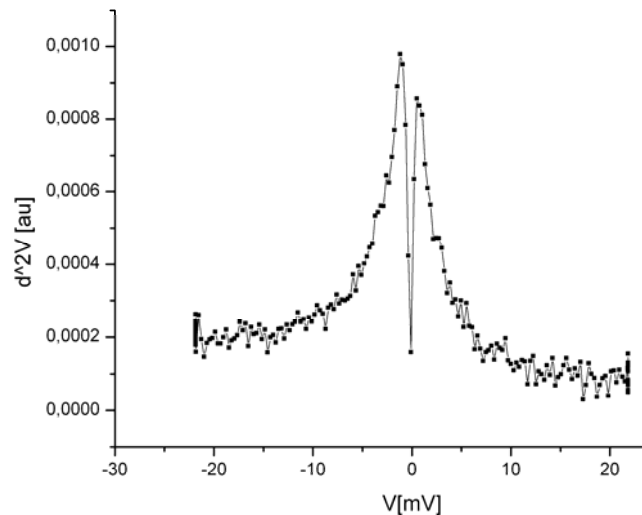
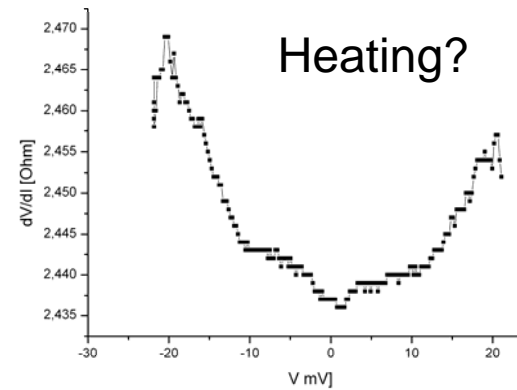
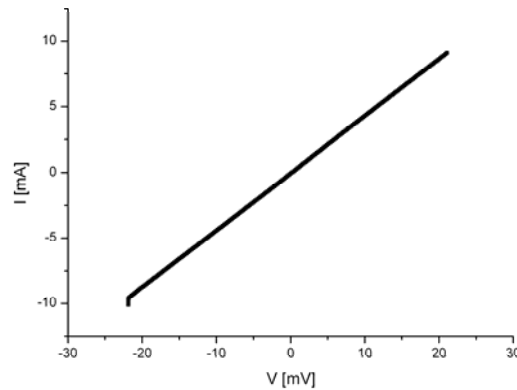


Fig. 1. Point-contact spectra of Cu single crystal in the orientation  $[100]$  (contact resistance  $R_0 = 2.75 \Omega$ ) and  $[111]$  ( $R_0 = 4.45 \Omega$ ). Temperature  $T = 1.6 \text{ K}$ ; the modulation amplitude  $V_1 = 0.2 \text{ mV}$  in the first case and  $V_1 = 0.3 \text{ mV}$  in the second case. The lower curve shows the expected energy dependence of the background of the PC spectrum.

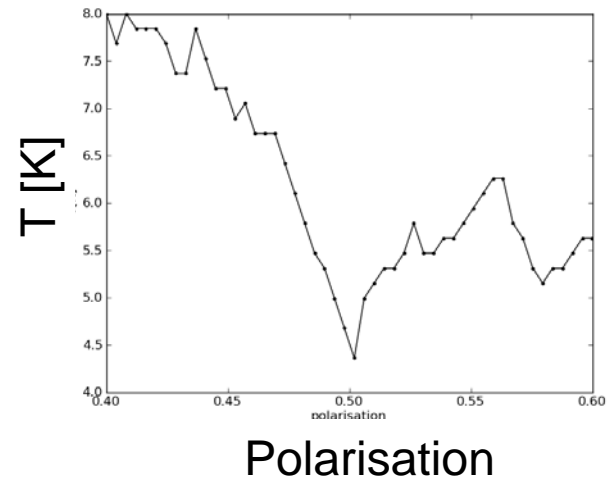
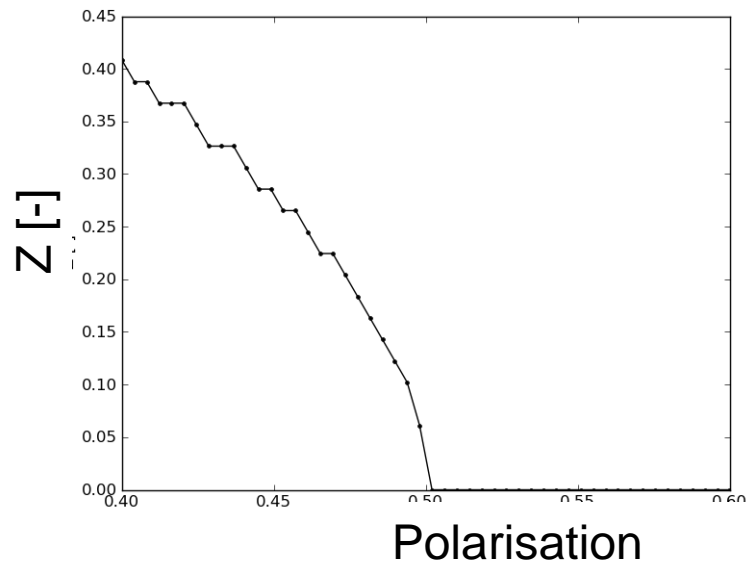
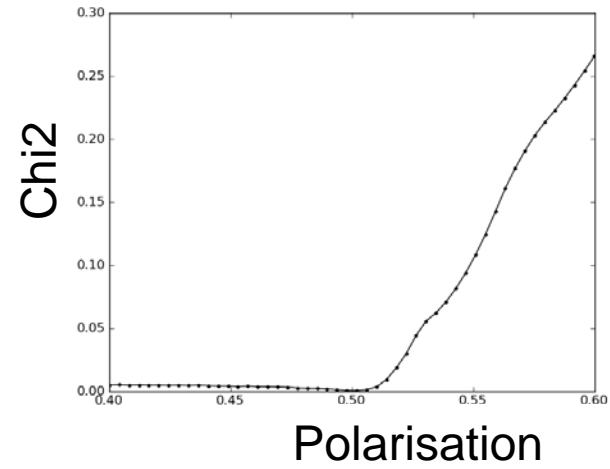
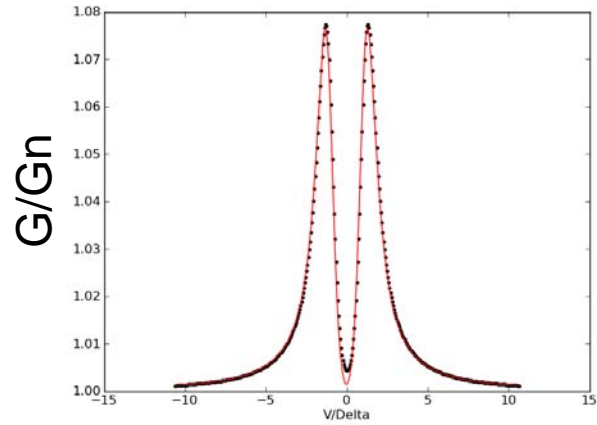
# Mechanical Point contacts



# Fitting

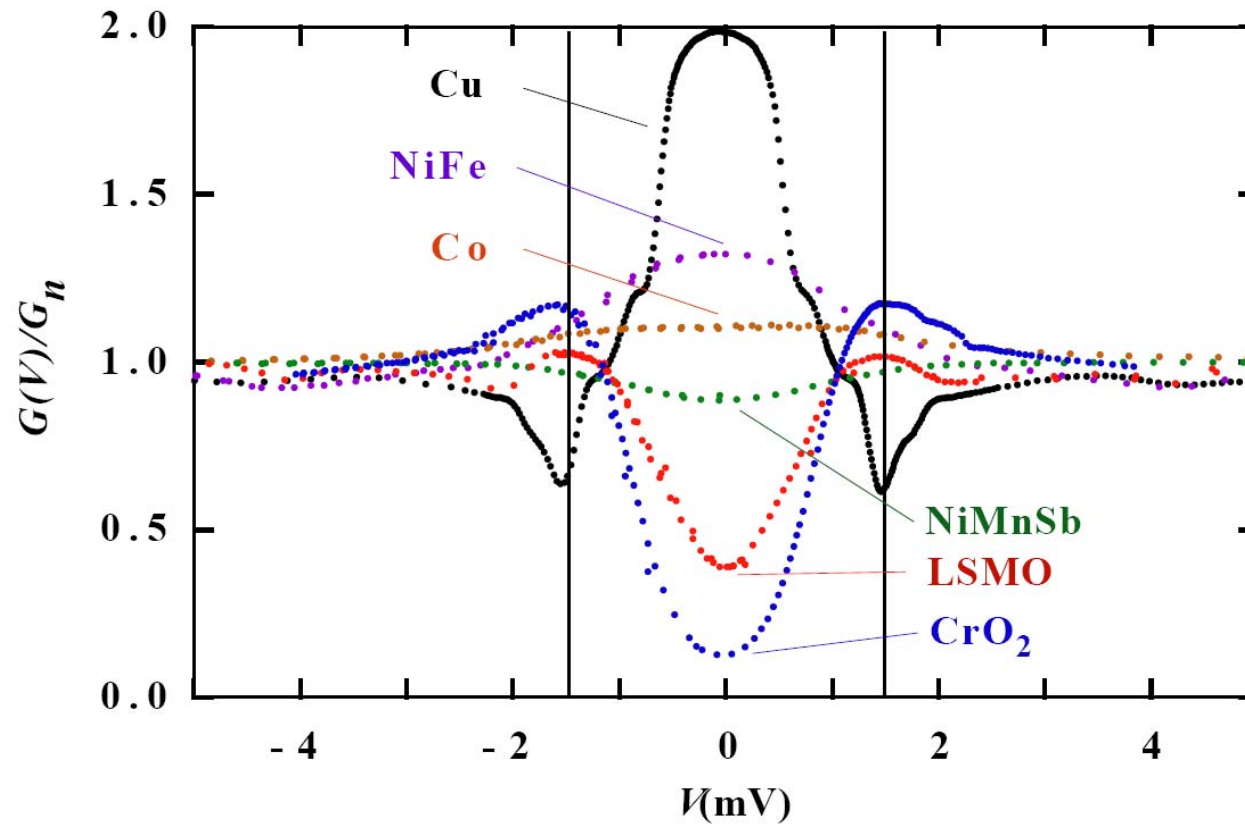
- $\chi^2 = \sum [G_{meas} - G_{fit}]^2$
- 3 parameters (Z, T, gap) -> unique solution
- 4 parameters (Z, T, gap, P) -> no unique solution
- How to normalize?

# Fitting



# Add spin

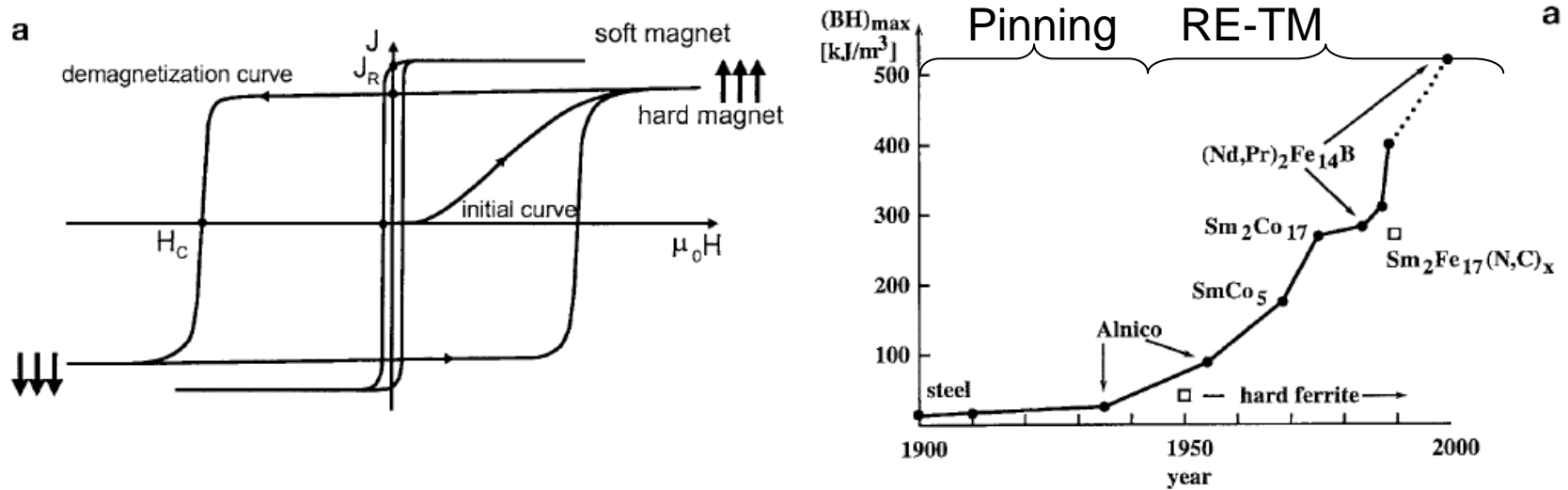
$$G = [1 - P]G_u(\Delta, Z, \omega) + PG_p(\Delta, Z, \omega)$$



- Spin flip laser
- Point contact spectroscopy
- **SmCo5**
- Radiation setup



# Hard magnets



Rare earth transition metals

- |      |      |
|------|------|
| - Nd | - Fe |
| - Pr | - Co |
| - Sm |      |

## A Family of New Cobalt-Base Permanent Magnet Materials

K. STERNAT, G. HOFFER, J. OLSON, AND W. OSTERTAG

*Air Force Materials Laboratory, Dayton, Ohio*

AND

J. J. BECKER

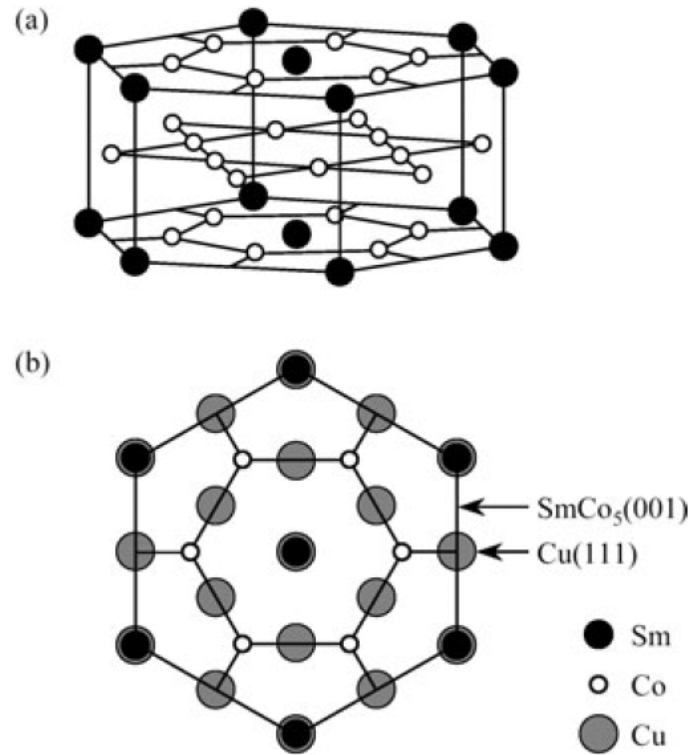
*General Electric Research and Development Center, Schenectady, New York*

The magnetocrystalline anisotropy of several intermetallic phases of the type  $\text{RCo}_5$  ( $\text{R} = \text{Y}, \text{Ce}, \text{Pr}, \text{Sm}$ , Y-rich and Ce-rich mischmetals) has been investigated, and it is concluded that these alloys are promising candidates for fine-particle permanent magnets. They have extremely high uniaxial anisotropy ( $K = 5.4$  to  $7.7 \times 10^7$  erg/cm<sup>3</sup>), single easy axis, high saturation ( $B_s = 8500$  to  $11\,200$  G) and Curie point ( $t_c = 464^\circ$  to  $747^\circ\text{C}$ ). Approximate upper limits for the possible energy product lie between 18 and 31.3 MGOe. Experimentally, coercive forces of over 8000 Oe and  $(BH)_{\text{max}} = 5.1$  MGOe have been observed in  $\text{SmCo}_5$  merely ground at room temperature. Grinding of  $\text{YCo}_5$  and  $(\text{Ce-MM})\text{Co}_5$  produces an increase of  $MH_c$  to 2200 and 2700 Oe, respectively, followed by a decrease as particle size continues to decrease.

### Combine

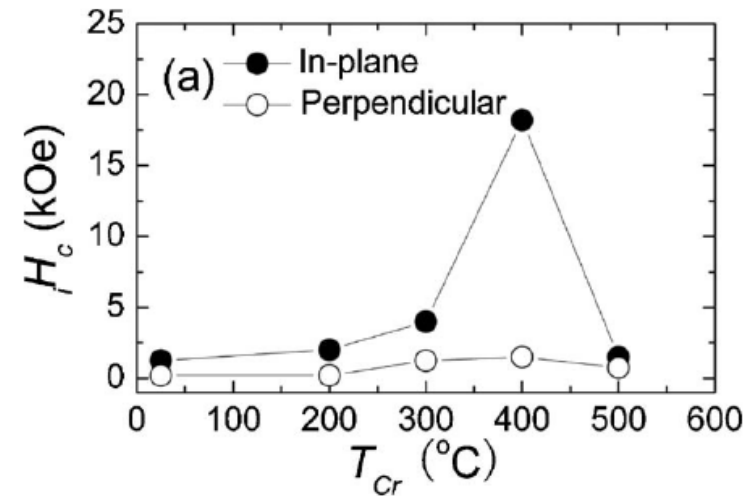
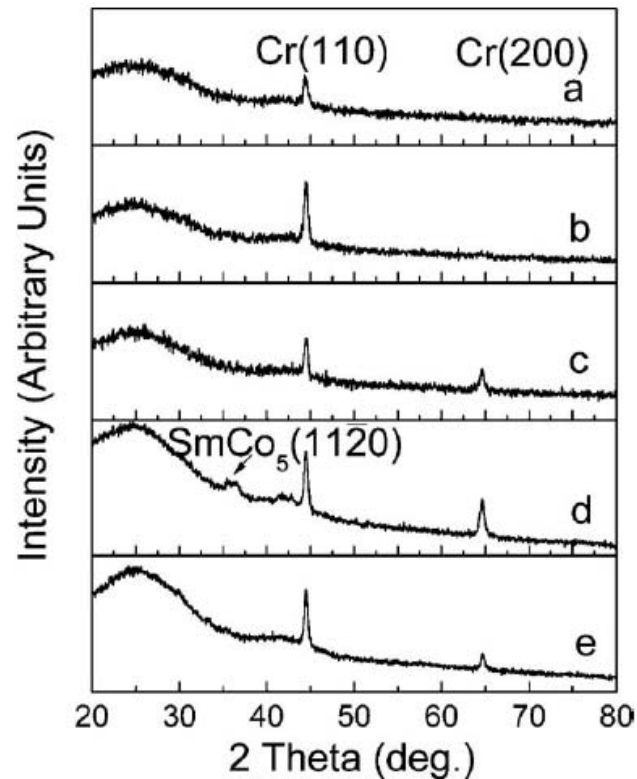
- High saturation polarization and  $T_{\text{curie}}$  3d TM
- high crystal anisotropy RE

# SmCo<sub>5</sub>



P6/mmm-structure

# Growing SmCo5



\* Claims  $H_c$  can be >6T @10K

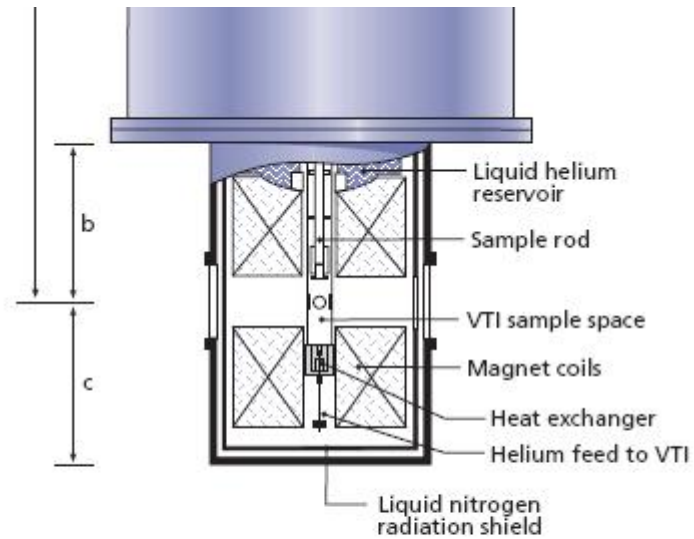
FIG. 2. The XRD spectra of Cr/SmCo/Cr thin films with the Cr underlayers deposited at different temperatures: (a) room temperature, (b) 200 °C, (c) 300 °C, (d) 400 °C, and (e) 500 °C.

Zhang et al. J.Appl Phys **103** 113908

\*Journal of Physics: Conference Series **234** (2010) 012012

- Spin flip laser
- Point contact spectroscopy
- SmCo5
- **Radiation setup**

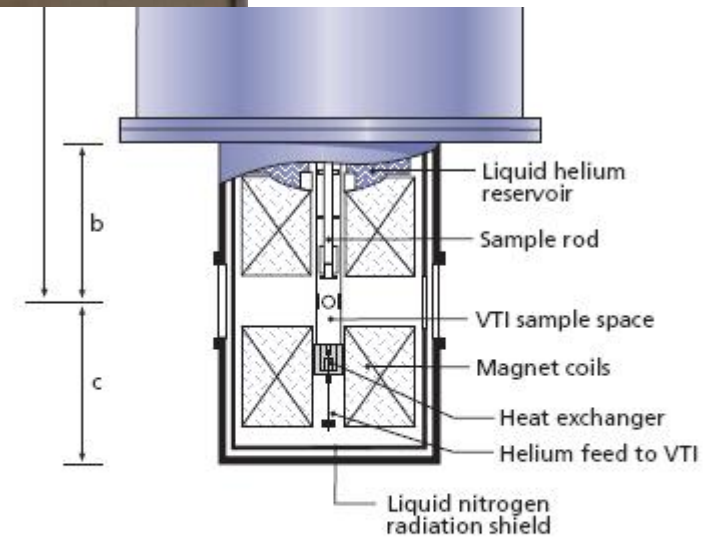
# Build setup to irradiate PC



Basic spectromag SM4000 system configuration  
Dimensions a, b and c are specified in the  
system configuration table.



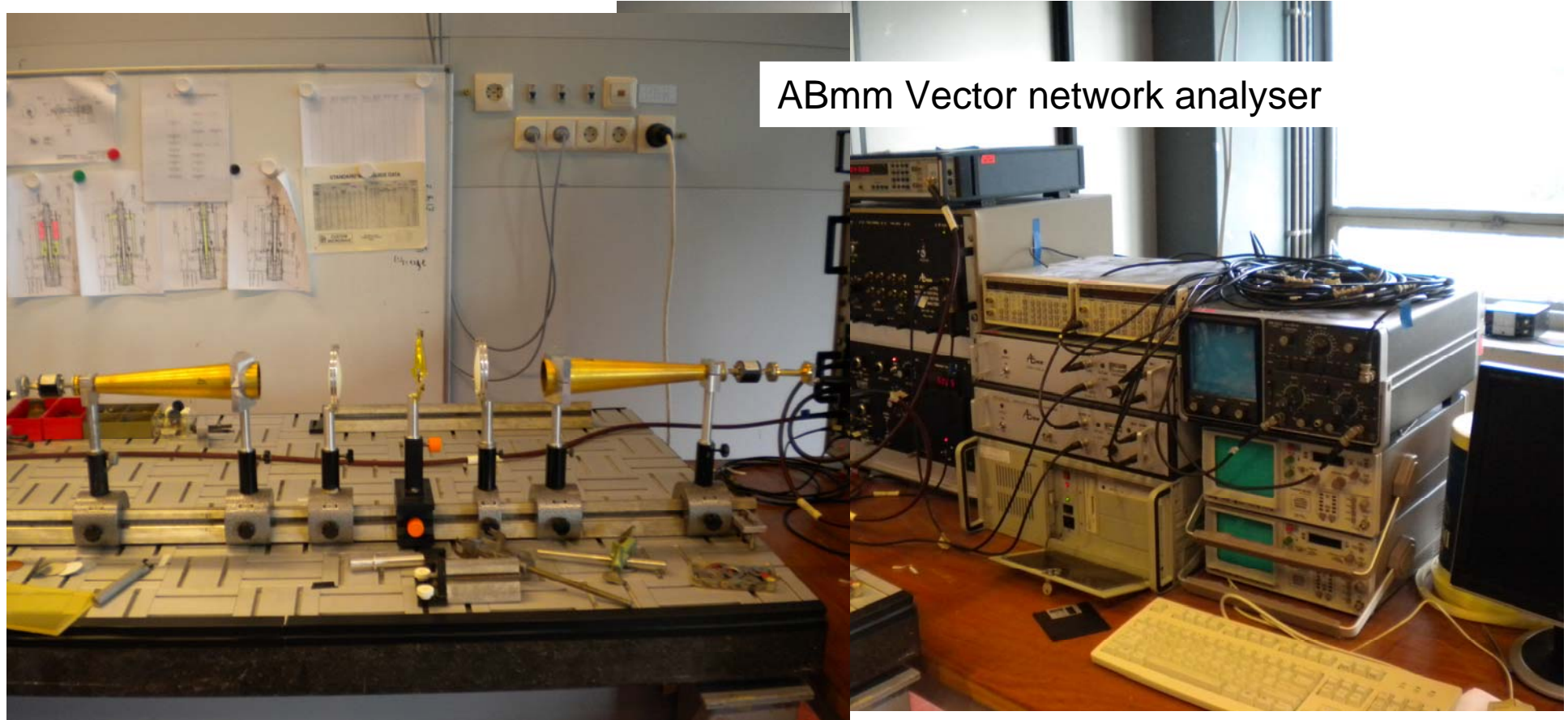
# Windows



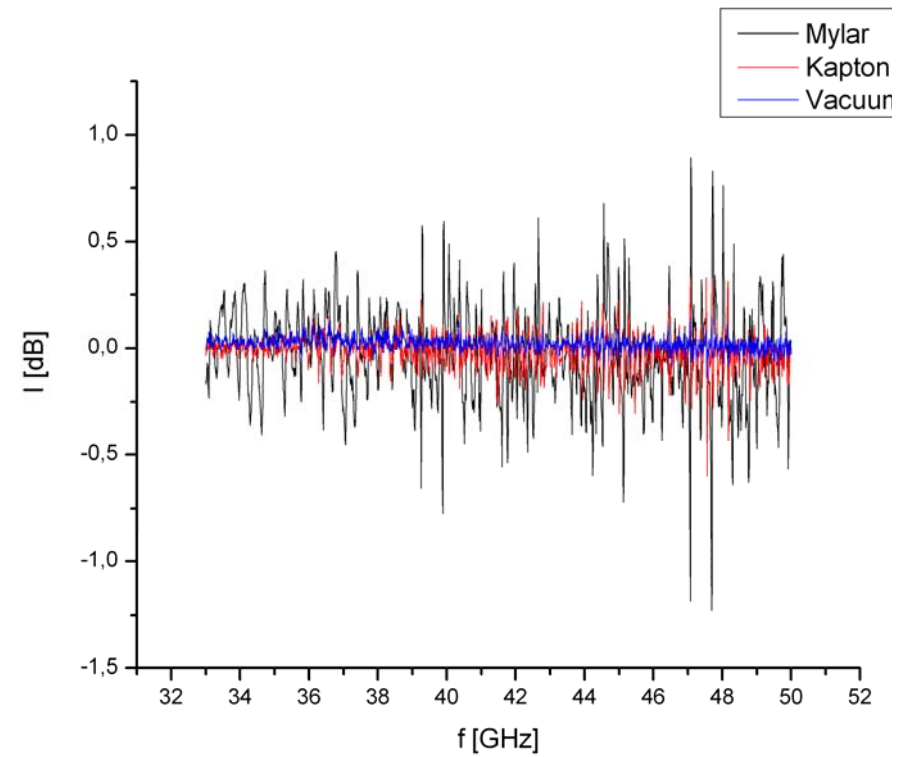
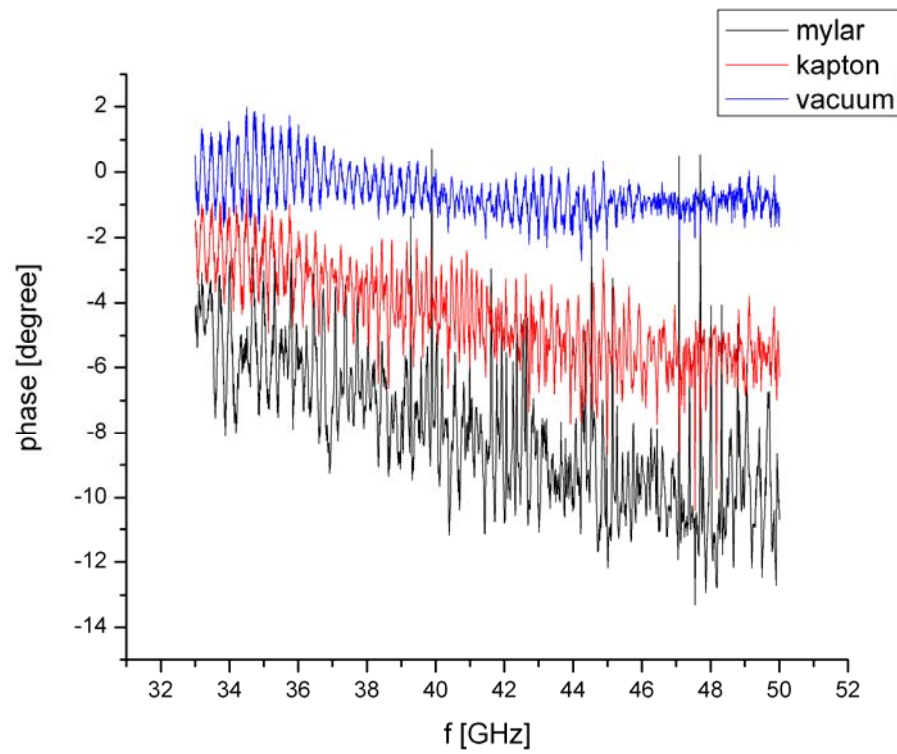
Basic spectromag SM4000 system configuration  
Dimensions a, b and c are specified in the  
system configuration table.



# Add radiation / Transmission



# Transmission window



# Outlook / To Do

- Study possible materials using PCS
- Grow SmCo<sub>5</sub>
- Measure PCS with radiation